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Thermodynamic and acoustic [properties](http://www.elsevier.com/locate/tca) [of](http://www.elsevier.com/locate/tca) [binary](http://www.elsevier.com/locate/tca) [m](http://www.elsevier.com/locate/tca)ixtures of oxolane with aniline and substituted anilines at 303.15, 313.15 and 323.15 K

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abstract

Speeds of sound u, isentropic compressibilities κ_{S} , Rao's molar sound functions R, intermolecular free lengths L_f , specific acoustic impedances Z, excess molar volumes V_m^E , excess isentropic compressibilities κ^E_S , excess intermolecular free lengths L_f^E and excess specific acoustic impedances Z^E , of three binary mixtures of oxolane (tetrahydrofuran) with aniline, N-methylaniline and N-ethylaniline have been reported over the entire range of composition at 303.15, 313.15 and 323.15 K. The excess values have been fitted to Redlich–Kister polynomial equation. The results have been analyzed in terms of molecular interactions between oxolane and anilines. The speeds of sound in present binary mixtures have been estimated from various empirical and theoretical models.

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1. Introduction

The thermodynamic, acoustic and transport properties of nonelectrolyte liquid–liquid mixtures provide information about type and extent of molecular interactions, and can be used for the development of molecular models for describing the behaviour of solutions [1–6]. They are also necessary for engineering calculation, research of mass transfer, heat transfer and fluid flow. The increasing use of cyclic ethers and alkyl and aromatic amines in many industrial processes, as well as theoretical interest in the nature of associated solutions have greatly stimulated the need for [exten](#page-7-0)sive information on properties of mixtures involving these components. It has been reported that cyclic ethers interact with amines in their mixtures. Different volumetric and thermal effects are observed upon mixing because of the molecular interactions between cyclic ethers + amine [7–13]. The formation of hydrogen bonds is assumed to occur between a primary or secondary amine group with weak proton donor ability and the unshared electron pairs on the oxygen atom of cyclic ether. Considerable systematic work on cyclic ether + alkylamine has been reporte[d,](#page-7-0) whereas data on cyc[lic](#page-7-0) [ether](#page-7-0) + aromatic amines is scanty [14,15]. Considering all these aspects we undertook investigations on the

thermodynamic, acoustic and transport properties of binary mixtures involving oxolane and aromatic amines. Oxolane is used as a solvent in many chemical industries due to its proton accepting nature. The aniline is predominantly used [16] as parent substance in the manufacture of several chemical products and intermediates. It is also used in manufacture of synthetic dyes, drugs and as an accelerator in vulcanization of rubber. Secondary amines N-methylaniline and N-ethylaniline are used as a latent and coupling solvent. In the present p[aper,](#page-7-0) [w](#page-7-0)e report densities, speeds of sound, isentropic compressibilities, Rao's molar sound functions, intermolecular free lengths, specific acoustic impedances, and calculated excess functions of three binary mixtures of oxolane with aniline, N-methylaniline, and N-ethylaniline at 303.15, 313.15 and 323.15 K. This study will also provide a test of Nomoto's relation (NR) [17], Van Dael and Vangeel relation (VVR) [18], Junjie's relation (JR)[19], impedance relation (IR)[20], Schhaffs' collision factor theory (CFT) [21,22], Jacobson's free length theory (FLT) [23], and Progogine–Flory–Patterson–Oswal theory (PFPOT) [24–27] to estimate speed of sound in binary mixtu[res at](#page-7-0) different temperatures.

2. [Experim](#page-7-0)ental

All chemicals used in this study were of analytical grade and obtained from s.d.fine-chem., Ltd. The claimed mass fraction purity for the chemicals was >0.995. These liquids were dried over 4\AA molecular sieves and partially degassed prior to use. The purity of

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Table 1 Densities (ρ) and speeds of sound (*u*), thermal expansion coefficients (α), and heat capacities (C_P) for pure liquids.

Liquid	T(K)	ρ (kg m ⁻³)			$u (m s^{-1})$		C_P (J mol ⁻¹ K ⁻¹)
		Expt.	Lit.	Expt.	Lit.		
Oxolane	303.15	877.3	877.0 ^a 876.7 ^d	1255.9	1256.3 ^b 1255 ^d	1.248	125 ^c
	313.15	867.2	865.4c 870.5f	1211.7	1209 ^e 1206.4 ^c	1.269	128 ^c
	323.15	856.6		1170.9		1.291	130 ^c
Aniline	303.15	1012.8	1012.9 ^g 1012.8 ^h	1615.2	1614.5^{i}	0.851^{j}	191 ^k
	313.15	1004.9	1004.4 ^h $1004.6^{\rm m}$	1582.6	1574.4^{i}	0.858^{j}	194 ¹
	323.15	998.4		1558.2		0.866^{j}	197 ¹
N-methyl aniline	303.15	978.2	$977.71^{\rm m}$ 978.1 ^h	1548.3	$\overline{}$	0.816^{j}	211 ¹
	313.15	969.8	969.9 ^h 969.6 ^m	1512.4		0.822^{j}	214 ¹
	323.15	961.0	961.9 ^h	1477.3		0.830^{j}	218 ¹
N-ethyl aniline	303.15 313.15 323.15	951.9 944.3 935.9	952.6 ^m 944.2 ^m	1497.4 1462.2 1430.5	$\overline{}$	0.885^{j} 0.893 0.902^{j}	227 ¹ 230 ¹ 233 ¹

^a Ref. [28].

^b Ref. [29].

 c Ref. [30].

^d Ref. [31]. ^e Ref. [32].

^f Ref. [33].

^g Ref. [34].

^h Ref. [35].

ⁱ Ref. [36].

^j Derived from density data from Ref. [37].

^k Ref. [38].

¹ Estimated from group contribution method of Chueh-Swanson [39].

^m Ref. [37].

these experimental li[quids](#page-7-0) was checked by comparing the observed [d](#page-7-0)ensities and velocities with those reported in the literature. The [m](#page-7-0)easured values are included in Table 1 a[long](#page-7-0) with the available literature values [28–38]. The agreement between the two is with in combined experimental error.

Airtight stoppered bottles were used for the preparation of the mixtures. The weight of the dry bottle was first determined. The less volatile component of the binary mixtures was introduced first in the sa[mple](#page-7-0) [bott](#page-7-0)le followed by second component, and the weight at each step was taken using an electronic balance (Mettler-AE 240, Switzerland) accurate to ± 0.05 mg. The densities of pure liquids and the binary mixtures were measured by calibrated single stem capillary pycnometer. The speeds of sound in pure liquids and in their binary mixtures were measured using single-crystal variable-path ultrasonic interferometer (Mittal Enterprises, New Delhi, Model: M-82) operating at 2 MHz. For all the measurements, temperatures were controlled by circulating the water through an ultra thermostat JULABO F-25 (made in Germany) keeping temperature fluctuations within ± 0.02 K. The details of the experimental procedure have been described elsewhere [40]. The uncertainty in density and speed of sound measurements was within 0.1 kg m⁻³ and 1 m s^{-1} .

3. Results and discussion

The results for the densities ρ , speeds of sound u, isentropic compressibilities κ_S , Rao's molar sound functions R [41], specific acoustic impedances Z and intermolecular free lengths L_f , excess molar volumes V_m^E and excess isentropic compressibilities κ_S^E for binary mixtures of oxolane with aniline, N-methylaniline and Nethylaniline at 303.15, 313.15, and 323.1[5 K ove](#page-7-0)r the entire range of composition are given in Tables 2–4.

From the values of densities and speeds of sound, the isentropic compressibilities, Rao's molar sound functions, specific acoustic impedances and intermolecular free length were obtained using the relations

$$
\kappa_{s} = \left(\rho u^{2}\right)^{-1} \tag{1}
$$

$$
R = u^{1/3}V
$$
 (2)

$$
Z = u\rho \tag{3}
$$

$$
L_f = \frac{K_{Jac}}{u\rho^{1/2}} = K_{Jac} \kappa_S^{-1/2}
$$
\n(4)

where $V = \sum (x_i M_i)/\rho$ in which x_i and M_i are the mole fraction and molecular mass of component *i.* $K_{lac} = ((91.368 + 0.3565T)10^{-8})$ is temperature dependent Jacobson's constant [23]. The values of $\kappa_{\mathsf{S}},$ R, and Z were uncertain within ± 0.2 % while L_f within 1%.

Excess molar volume V_m^E , excess isentropic compressibility κ_S^E , excess intermolecular free length L_f^E and excess specific acoustic impedance Z^E in each mixture [were ca](#page-7-0)lculated from ρ and κ_S of pure liquids and binary mixtures with following expression

$$
Y^E = Y - Y^{id} \tag{5}
$$

where Y represent either V, κ_S , Z and L_f . The V $^{id}_m$ for an ideal mixture was calculated from usual relation

$$
V_m^{id} = \Sigma x_i V_i^o = \frac{\Sigma x_i M_i}{\rho_i} \tag{6}
$$

 $\kappa^{\sf id}_{\mathsf{S}}$ for an ideal mixture was calculated from the relation recommended by Benson and Kiyohara [42] and Douheret et al. [43]

$$
\kappa_s^{id} = \Sigma \phi_i \left\{ \frac{\kappa_{S,i}^o + TV_i^o(\alpha_i^o)^2}{C_{p,i}^o} \right\} - T(\Sigma x_i V_i^o) \left(\frac{\Sigma \phi_i \alpha_i^{o2}}{\Sigma x_i C_{p,i}^o} \right) \tag{7}
$$

Table 2 Densities (ρ), speeds of sound (u), isentropic compressibilities (κ_S), Rao's molar sound functions (R), specific acoustic impedances (Z) and intermolecular free lengths (L_f), excess molar volumes (V_m^E) and excess isentropic compressibilities (κ_S^E) for oxolane (1)+aniline (2) mixture at 303.15, 313.15 and 323.15 K.

 L_f^{id} and Z^{id} for an ideal mixture were obtained as

$$
L_f^{id} = K_{Jac}(\kappa_S^{id})^{1/2}
$$
\n
$$
Z^{id} = \rho^{id} u^{id} = \left(\frac{\rho^{id}}{\kappa_S^{id}}\right)^{1/2}
$$
\n(9)

in which the V_i^0 , α_i^0 and $C_{p,i}^0$ are, respectively, the molar volume, isobaric thermal expansion coefficient and molar isobaric heat capacity of pure component i, $\phi_i = x_iV_i/\sum x_iV_j$ is the volume fraction of i in the mixture, stated in terms of the unmixed components and $\rho^{id} = \Sigma \phi_i \rho_i^o$. The values of α^o and C_ρ^o used for these calculations are listed in Table 1.

For compact and smooth representation, the values of u and $\kappa_{\mathcal{S}}$ were fitted to a polynomial equation of the form

$$
u \text{ or } \kappa_S = \sum_{i=1}^{m} A_i X_1^{i-1} \tag{10}
$$

The excess molar volume V_m^E , excess isentropic compressibility κ^E_{S} , excess intermolecular free length L^E_{F} and excess specific acoustic impedance Z^E were correlated by Redlich–Kister polynomial equation [44]

$$
Y^{E} = x_{1}(1 - x_{1}) \sum_{i}^{n} B_{i}(1 - 2x_{1})^{i}
$$
 (11)

where $Y^E = V^E_m$, κ^E_S , L^E_f and Z^E .

The coefficients A_i of Eq. (10) and B_i of Eq. (11) and corresponding standard deviations σ obtained from a least-squares fit with equal weights assigned to each point are given in Tables 5 and 6.

It is observed from Tables 2–4, that the values of speeds of sound u, Rao's molar sound functions R, specific acoustic impedances Z decrease with mole fraction of oxolane and also with rise in temperature for all the systems studied. The values of $\kappa_{\mathcal{S}}$ and L_f increase with the mole fraction of oxola[ne](#page-5-0) [and](#page-5-0) [with](#page-5-0) [the](#page-5-0) rise in temperature. The Figs. 1–4 shows the dependence of V_m^E, κ_S^E, L_f^E , and Z^E on composition. V_m^E and κ_S^E are negative for all the systems over the entire mole fraction range and their magnitude decreases with rise in temperature from 303.15 to 323.15 K. The minimum κ_S^E occurs [at 0.6 mo](#page-3-0)l fraction for all the systems.

For oxolane + aniline, speed of sound and excess molar volumes at 308.15 K have been reported by Deshpande et al. [15] and excess molar volumes at 298.15 K by Suri and Naorem [14]. The present value of equimolar excess molar volume for oxolane + aniline at 303.15 K is 0.896×10^{-6} m³ mol⁻¹ which is in between the values -0.916×10^{-6} m³ mol⁻¹ of Suri and Naorem [14] at 298.15 K and -0.875×10^{-6} m³ mol⁻¹ of Deshpa[nde](#page-7-0) [et](#page-7-0) al. [15] at 308.15 K. Considering the effect of temperature [the](#page-7-0) [ag](#page-7-0)reement is very good. Our inter-polated values of excess isentropic compressibility of equimolar oxolane + aniline mixture at 308.15 K is -80.2 TPa⁻¹ while Deshpande et al. [15] reported -63 -63 [T](#page-7-0)[Pa](#page-7-0)⁻¹. For other two mixtures no work on either speed of soun[d](#page-7-0) [or](#page-7-0) [ex](#page-7-0)cess molar volumes at any temperature seems to have been reported in the literature.

The overall behaviour of V_m^E and κ_S^E can be envisaged as a resultant of opposite effects: (i) disruption of ether–ether interaction[s,](#page-7-0) [\(ii\)](#page-7-0) break down of dipolar and/or H-bonding self-

Fig. 1. Dependence of excess molar volumes on mole fraction of oxolane, (a) oxolane + aniline, (b) oxolane + N-methylaniline and (c) oxolane + N-ethylaniline at $T = 303.15$ K (\bullet), 303.15 K (\blacksquare), and 323.15 K (\blacktriangle).

association in oxolane and anilines, and (iii) aniline–oxolane cross-association and dipole–dipole interactions. The negative values of V_m^E and κ_S^E suggest qualitatively that dipole-dipole interactions (μ for oxolane = 1.75 D and for anilines \approx 1.51 D [38]) and cross-association between aniline and oxolane components dominate over the dispersive ether–ether interactions, breaking of dipolar order of aniline as well as of oxolane and the disruption of H-bonding in primary (aniline) and secondary amines (N-methyl and N-ethylanilines). In the cross-association t[he](#page-7-0) [H-a](#page-7-0)tom of the aniline and substituted anilines interact with O-atom of the oxolane molecules. The negative κ_S^E and V_m^E values at equimolar concentra-

Fig. 2. Dependence of excess isentropic compressibilities on mole fraction of oxolane, (a) oxolane + aniline, (b) oxolane + N-methylaniline and (c) oxolane + Nethylaniline at T = 303.15 K (●),303.15 K (■), and 323.15 K (▲).

tions follow the order aniline > N-ethylaniline > N-methylaniline. Apparently, from the values of excess functions, one can say that the strength of specific interactions between unlike molecules varies in the order aniline > N-methylaniline > N-ethylaniline. The decreased interactions from aniline to N-substituted aniline are caused by steric hindrance due to alkyl group attached to N-atom and strength of base. This observation is in accordance with the findings of Fort and Moore [45]. The decrease in the magnitude of negative κ_S^E and V_m^E with the rise in temperature further suggests that interactions between unlike components tend to reduce due to the increase in thermal motions.

As expected, the trend of negative L_f^E values (Fig. 3 at 303.15 K) is similar to κ_S^E in all the systems. The negative values of L_f^E are generally observed in systems with specific interactions between unlike molecules. Fig. 4 shows that Z^E is positive for all the systems. Specific acoustic impedance is a quantity, which depends on the molecular packing of the systems. The positive values of Z^E are in accordance with the presence of strong interactions between the

unlike molecules. Positive trends in Z^E have also been reported for binary mixtures of methanol + piperidine [46].

The dependence of Rao's molar sound functions R on mole fractions x_1 is shown in Fig. 5. The Rao's molar sound functions R decrease with mole fraction of oxolane and are independent of temperature (Tables 2–4). The plots show that there are slight negative deviations from the linear de[pende](#page-7-0)nce on x_1 for all three binary mixtures a[t all tem](#page-6-0)peratures. According to Rao's approach [41], the

Fig. 3. Dependence of excess intermolecular free length on mole fraction of oxolane at T=303.15 K (a) oxolane + aniline (\bullet), (b) oxolane + N-methylaniline (\blacksquare) and (c) oxolane + N-ethylaniline (\triangle) .

Fig. 4. Dependence of excess specific acoustic impedance on mole fraction of oxolane at T=303.15 K (a) oxolane + aniline (\bullet), (b) oxolane + N-methylaniline (\blacksquare) and (c) oxolane + N-ethylaniline (\triangle) .

Table 4 Densities (ρ), speeds of sound (u), isentropic compressibilities (κ_S), Rao's molar sound functions (R), specific acoustic impedances (Z) and intermolecular free lengths (L_f), excess molar volumes (V_m^E) and excess isentropic compressibilities (κ_S^E) for oxolane (1)+N-ethylaniline (2) mixture at 303.15, 313.15 and 323.15 K.

speed of sound is directly related to the strength of intermolecular energy. If the intermolecular energy were simply a consequence of additive non-bonding interaction, the Rao's molar sound function should predict additive behaviour. The fact that is not so implies some additional interactions in the mixtures compared with those in the simple liquids [47].

4. Estimation of speed of sound

The speed of sound u from the Nomoto's relation [17], Van Dael and Vangeel [18], Junjie's relation [19], impedance relation [20], collision factor theory [21,22], intermolecular free length theory [23], and Prigogine–Flory–Patterson–Oswal theory [24–27] have

Table 5 Coefficients A_i of Eq. (10) along with standard deviations σ of binary mixture properties.

Properties	Temperature	A ₀	A ₁	A ₂	A_3		$\sigma(Y^E)$		
Oxolane (1) + aniline (2)									
$u (m s^{-1})$	303.15	1615.4	-278.3	4.9	-189.9	103.7	1.1		
	313.15	1584.0	-289.8	-35.3	-47.4		1.5		
	323.15	1558.8	-246.9	-205.1	203.8	-139.8	1.5		
κ_S (TPa ⁻¹)	303.15	378.5	165.7	45.6	133.4		0.9		
	313.15	396.8	166.4	175.6	-144.5	191.6	1.0		
	323.15	412.4	142.2	359.1	-490.3	427.8	1.0		
Oxolane (1) + N-methylaniline (2)									
$u (m s^{-1})$	303.15	1548.4	-258.4	101.9	-279.5	143.6	0.4		
	313.15	1511.9	-225.8	-76.0	1.0		0.8		
	323.15	1477.6	-225.1	-78.7	-3.1		0.4		
κ_S (TPa ⁻¹)	303.15	427.0	153.2	59.0	84.2		0.8		
	313.15	450.7	174.4	32.7	127.6		0.3		
	323.15	476.0	189.9	25.1	160.0		0.7		
Oxolane (1) + N-ethylaniline (2)									
$u (m s^{-1})$	303.15	1497.1	-152.1	38.4	-253.0	125.4	0.8		
	313.15	1461.8	-115.5	-130.7	-4.4		0.6		
	323.15	1430.6	-106.6	-164.8	12.2		0.6		
κ_S (TPa ⁻¹)	303.15	469.2	104.0	43.2	107.1		0.8		
	313.15	495.1	116.5	20.0	153.7		0.4		
	323.15	522.4	82.5	219.6	-130.1	156.2	0.7		

also been estimated for the present binary mixtures. The pertinent relations in these calculations and their theoretical basis have been outlined several times and will not be repeated here.

The speeds of sound u_{est} in the binary mixture over the entire range of composition were estimated. The standard percentage deviation σ % was calculated by considering dev% for all mole fractions using the relation

$$
\sigma \mathscr{E} = \left[\frac{\Sigma (dev \mathscr{E})^2}{n-1} \right]^{1/2} \tag{12}
$$

where *n* represents number of data points.

The study of Table 7 reveals that the values of σ % for the NR, VVR, JR, IR, CFT, FLT, and PFPOT are in the range from 0.1 to 2.1, 5.1 to 8.0, 1.4 to 5.1, 0.4 to 1.9, 0.4 to 1.8, 1.2 to 2.8 and 1.0 to 3.3, respectively for the three binary mixtures at three temperatures investigated. The average values of standard deviations $\langle \sigma \mathcal{X} \rangle$ for the NR, IR and CFT are comparatively less than those found for other estimations.

Fig. 5. Dependence of Rao's molar sound function on mole fraction of oxolane at T=303.15 K (a) oxolane + aniline (\bullet), (b) oxolane + N-methylaniline (\blacksquare) and (c) oxolane + N-ethylaniline (\triangle) .

Table 7

Standard percentage deviation (σ %) between experimental and estimated speed of sound using empirical or theoretical relations for binary mixtures.

Thus, the average values of standard deviation, the estimation ability of speed of sound in the presently investigated mixtures follows the sequence CFT \cong NR > IR > FLT > PFPOT > JR > VVD.

5. Conclusions

The values of V_{m}^{E},κ_{S}^{E} and L_{f}^{E} are negative and Z^{E} are positive for binary mixtures of oxolane with aniline, N-methylaniline and Nethylaniline at 303.15, 313.15, and 323.15 K. There exist specific interactions between unlike molecules through hydrogen bonding and dipole–dipole interactions between unlike molecules. The estimation ability of speed of sound in the presently investigated mixtures follows the sequence CFT ≅ NR > IR > FLT > PFPOT > JR > VVD.

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